Deformation Problem in Micropolar Generalized Thermoelastic Medium using various Transform Techniques

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Abstract: The present investigation deals with the deformation in micropolar generalized thermoelastic medium with mass diffusion subjected to thermo mechanical loading due to thermal laser pulse. Laplace and Fourier transform technique is used to solve the problem. Concentrated normal force and thermal source are taken to illustrate the utility of approach. The closed form expressions of normal stress, tangential stress, tangential couple stress, mass concentration and temperature distribution are obtained in the transformed domain. Numerical inversion technique of Laplace transform and Fourier transform has been applied to obtain the resulting quantities in the physical domain after developing a computer program. The normal stress, tangential stress, tangential coupled stress, temperature distribution and mass concentration are depicted graphically to show the effect of relaxation times. Some particular cases of interest are deduced from the present investigation.

Keywords: Laser, Pulse, Micropolar, Mass diffusion, uniformly and linearly distributed source.

Introduction

Micropolar theory of elasticity was introduced by Eringen [1]. This theory incorporates the local deformation and rotation of the material points of the composite. This theory provides a model that can support body couples and surface couples and exhibits a high frequency optical wave spectrum. Eringen [2, 3], Maugin and Mild [4], Nowacki [5] developed the linear theory of micropolar thermoelasticity by excluding the micropolar theory of elasticity to include the thermal effects. Touchert et al. [6] derived the basic equations of linear theory of micropolar coupled thermoelasticity. Diffusion is defined as the spontaneous movement of the particles from a high concentration region to the low-concentration region, and it occurs in response to a concentration gradient expresses as the change in the concentration due to change in position. Thermal diffusion utilizes the transfer of heat across a thin liquid or gas to accomplish isotope separation. Today, thermal diffusion remains a practical process to separate isotopes of noble gases e.g., Xenon and other light isotopes e.g., Carbon for research purposes. In most of the applications, the concentration is calculated using Fick’s law. This into consideration is a simple law which does not take the mutual interaction between the introduced substance and the medium into which it is introduced or the effect of temperature of this interaction. However, there is a certain degree of coupling with temperature and temperature gradients as temperature speeds up the diffusion process. Nowacki [7, 10] developed the theory of thermoelastic diffusion by using coupled thermoelastic model. Dudziak and Kowalski [11] and Olesiak and Pyryev [12], respectively, discussed the theory of thermo diffusion and coupled quasi stationary problems of thermal diffusion for an elastic layer. Laser technology has a vital application in nondestructive materials testing and evaluation. When a solid is heated with a laser pulse, it absorbs some energy which results in an increase in localized temperature. This cause thermal expansion and generation of the ultrasonic waves in the material. There are generally two mechanisms for such wave generation, depending on the energy density deposited by the laser pulse. At high energy density, a thin surface layer of the solid material melts, followed by an ablation process whereby particles fly off the surface, thus giving rise to forces that generates ultrasonic waves. At low energy density, the surface material does not melt, but it expands at a high rate and wave and wave motion is generated due to thermoelastic processes. Very rapid thermal processes (e.g., the thermal shock due to exposure to an ultra-short laser pulse) are interesting from the stand point of thermoelasticity, since they require a coupled analysis of the temperature and deformation fields. A thermal shock induces very rapid movement in the structural elements, giving the rise to very significant inertial forces, and thereby, an increase in vibration. Rapidly oscillating contraction and expansion generates temperature changes in materials susceptible to diffusion of heat by conduction [13]. This mechanism has attracted considerable attention due to the extensive use of pulsed laser technologies in material processing and non-destructive testing and characterization [14, 15]. The so-called ultra-short lasers are those with pulse durations ranging from nanoseconds to femto seconds. In the case of ultra-short pulsed laser heating, the high intensity energy flux and ultra-short duration lead to a very large thermal gradients or ultra-high heating may exist at the boundaries. In such cases, as pointed out by many investigators, the classical Fourier model, which leads to an infinite propagation speed of the thermal energy, is no longer valid [16]. Researchers have proposed several models to describe the mechanism of heat conduction during short-pulse laser heating, such as the parabolic one-step model [17], the hyperbolic one-step model [18], and the parabolic two-step and hyperbolic two-step models [19, [20].

In this research, taking into account the mass concentration effect and radiation of ultra-short laser, we have established a model for micropolar thermoelastic medium with mass diffusion by using Laplace and Fourier transforms. The stress components and temperature distribution have been computed numerically. The resulting quantities are shown graphically to show the effect of mass concentration and temperature.

Problem Formulation

Following Eringen [3] and Al-Qahtani and Datta [30] the basic equations for homogeneous, isotropic micropolar generalized thermoelastic solid with mass diffusion in the absence of body forces and body couples are given by:

$$(\lambda+\mu)\nabla(\nabla.\boldsymbol{u})+(\mu+K)\nabla^2\boldsymbol{u}+K\nabla\times\boldsymbol{\phi}-\beta_1\left(1+\tau_1\frac{\partial}{\partial t}\right)\nabla T-\beta_2\left(1+\tau^1\frac{\partial}{\partial t}\right)\nabla C=\rho\ddot{\boldsymbol{u}},(1) \qquad (1)$$

$$(\gamma\nabla^2-2K)\boldsymbol{\phi}+(\alpha+\beta)\nabla(\nabla.\boldsymbol{\phi})+K\nabla\times\boldsymbol{u}=\rho j\ddot{\boldsymbol{\phi}}, (2)$$

$$K^*\nabla^2 T=\rho c^*\left(\frac{\partial}{\partial t}+\tau_0\frac{\partial^2}{\partial t^2}\right)T+\left(1+\varepsilon\tau_0\frac{\partial}{\partial t}\right)(\beta_1 T_0\nabla.\dot{\boldsymbol{u}}-Q)+aT_0\left(\frac{\partial}{\partial t}+\gamma_1\frac{\partial^2}{\partial t^2}\right)C, \qquad (3)$$

$$D\beta_2\nabla^2(\nabla.\boldsymbol{u})+Da\left(1+\tau_1\frac{\partial}{\partial t}\right)\nabla^2 T+\left(\frac{\partial}{\partial t}+\varepsilon\tau^0\frac{\partial^2}{\partial t^2}\right)C-Db\left(1+\tau^1\frac{\partial}{\partial t}\right)\nabla^2 C=0, \qquad (4)$$

$$t_{ij}=\lambda u_{k,k}\delta_{ij}+\mu(u_{i,j}+u_{j,i})+K(u_{j,i}-\epsilon_{ijk}\phi_k)-\beta_1\left(1+\tau_1\frac{\partial}{\partial t}\right)\delta_{ij}T-\beta_2\left(1+\tau^1\frac{\partial}{\partial t}\right)\delta_{ij}C, \qquad (5)$$

$$m_{ij}=\alpha\phi_{k,k}\delta_{ij}+\beta\phi_{i,j}+\gamma\phi_{j,i}, \qquad (6)$$

The plate surface is illuminated by laser pulse given by the heat input

$$Q=I_0 f(t)g(x_1)h(x_3) \qquad (7)$$

where I_0 is the energy absorbed. The temporal profile $f(t)$ is represented as,

$$f(t)=\frac{t}{t_0^2}e^{-\left(\frac{t}{t_0}\right)} \qquad (8)$$

Here t_0 is the pulse rise time. The pulse is also assumed to have a Gaussian spatial profile in x_1

$$g(x)=\frac{1}{2\pi r^2}e^{-\left(\frac{x_1^2}{r^2}\right)} \qquad (9)$$

Where r is the beam radius, and as a function of the depth x_3 the heat deposition dueto the laser pulse is assumed to decay exponentially within the solid,

$$h(x_3)=\gamma^* e^{-\gamma^* x_3} \qquad (10)$$

Equation (7) with the aid of (8,9 and 10) takes the form

$$Q=\frac{I_0\gamma^*}{2\pi r^2 t_0^2}te^{-\left(\frac{t}{t_0}\right)}e^{-\left(\frac{x_1^2}{r^2}\right)}e^{-\gamma^* x_3}\quad , \qquad (11)$$

we take

$$\boldsymbol{u}=(u_1,0,u_3),\boldsymbol{\phi}=(0,\phi_2,0), \qquad (12)$$

For further consideration, it is convenient to introduce in equations (1.1)-(1.4) the dimensionless quantities defined as:

$$u_i'=\frac{\rho\omega^* c_1}{\beta_1 T_0}u_i\ ,\quad x_i'=\frac{\omega^*}{c_1}x_i\ ,\quad t'=\omega^* t\ ,\quad T'=\frac{T}{T_0},\quad \tau_1'=\omega^*\tau_1\ ,\quad \tau_0'=\omega^*\tau_0\ ,\gamma_1'=\omega^*\gamma_1,\ t_{ij}'=\frac{1}{\beta_1 T_0}t_{ij}\ ,\omega^*=\frac{\rho c^* c_1^2}{K^*},\phi_i'=\frac{\rho c_1^2}{\beta_1 T_0}\phi_i\ ,\quad \tau^{1\prime}=\omega^*\tau^1\ ,\quad c_1^2=\frac{\lambda+2\mu+k}{\rho},\quad c_2^2=\frac{\mu+k}{\rho},\quad c_3^2=\frac{\gamma}{\rho j},\quad c_4^2=\frac{2\alpha_0}{\rho j_0},\quad \varepsilon=\frac{\gamma^2 T_0}{\rho^2 c^* c_1},m_{ij}^*=\frac{\omega^*}{c\beta_1 T_0}m_{ij}\ ,\quad C'=\frac{\beta_2}{\rho c_1^2}C\ ,$$

$$Q=\frac{K^*\omega^{*2}}{C^*}Q' \qquad (13)$$

Making use of (13) in (1)-(4) and with the aid of (12), we obtain:

$$a_1\frac{\partial e}{\partial x_1}+a_2\nabla^2 u_1-a_3\frac{\partial\phi_2}{\partial x_3}-\left(1+\tau_1\frac{\partial}{\partial t}\right)\frac{\partial T}{\partial x_1}-a_4\left(1+\tau^1\frac{\partial}{\partial t}\right)C=\rho\frac{\partial^2 u_1}{\partial t^2}, \qquad (14)$$

$$a_1\frac{\partial e}{\partial x_3}+a_2\nabla^2 u_3+a_3\frac{\partial\phi_2}{\partial x_1}-\left(1+\tau_1\frac{\partial}{\partial t}\right)\frac{\partial T}{\partial x_3}-a_4\left(1+\tau^1\frac{\partial}{\partial t}\right)C=\rho\frac{\partial^2 u_3}{\partial t^2}, \qquad (15)$$

$$\nabla^2\phi_2-2a_6\phi_2+a_6\left(\frac{\partial u_1}{\partial x_3}-\frac{\partial u_3}{\partial x_1}\right)=a_7\ddot{\phi}_2 \qquad (16)$$

$$-\nabla^2 T+\left(\frac{\partial}{\partial t}+\tau_0\frac{\partial^2}{\partial t^2}\right)T+a_5\left(\frac{\partial}{\partial t}+\varepsilon\tau_0\frac{\partial^2}{\partial t^2}\right)e+a_8\left(\frac{\partial}{\partial t}+\gamma_1\frac{\partial^2}{\partial t^2}\right)C=Q_0 f^*(x_1,t)e^{-\gamma^* x_3} \qquad (17)$$

$$\nabla^2 e+a_9\left(1+\tau_1\frac{\partial}{\partial t}\right)\nabla^2 T+a_{10}\left(1+\epsilon\tau^0\frac{\partial}{\partial t}\right)\dot{C}-a_{11}\left(1+\tau^1\frac{\partial}{\partial t}\right)\nabla^2 C=0\ , \qquad (18)$$

The displacement components u_1 and u_3 are related to the non- dimensional potential functions ϕ and ψ as:

$$u_1 = \frac{\partial\phi}{\partial x_1} - \frac{\partial\psi}{\partial x_3}\ ,\quad u_3 = \frac{\partial\phi}{\partial x_3} + \frac{\partial\psi}{\partial x_1} \qquad (19)$$

Substituting the values of $u_1$andu_3 from (19) in (14)-(18) and with the aid of (12), we obtain:

$$\nabla^2\phi - \ddot{\phi} - \left(1 + \tau_1\frac{\partial}{\partial t}\right)T - a_4\left(1 + \tau^1\frac{\partial}{\partial t}\right)C = 0, \qquad (20)$$

$$\nabla^4\phi + a_9\left(1 + \tau_1\frac{\partial}{\partial t}\right)\nabla^2 T + a_{10}\left(1 + \epsilon\tau^0\frac{\partial}{\partial t}\right)\dot{C} - a_{11}\left(1 + \tau^1\frac{\partial}{\partial t}\right)\nabla^2 C = 0, \qquad (21)$$

$$\left(1 + \tau_0\frac{\partial}{\partial t}\right)\dot{T} + a_5\left(\frac{\partial}{\partial t} + \varepsilon\tau_0\frac{\partial^2}{\partial t^2}\right)\nabla^2\phi + a_8\left(\frac{\partial}{\partial t} + \gamma_1\frac{\partial^2}{\partial t^2}\right)C - \nabla^2 T = Q_0 f^*(x_1,t)e^{-\gamma^* x_3}, \qquad (22)$$

$$a_2\nabla^2\psi - \ddot{\psi} + a_3\phi_2 = 0, \qquad (23)$$

$$\nabla^2\phi_2 - 2a_6\phi_2 - a_6\nabla^2\psi = a_7\ddot{\phi}_2, \qquad (24)$$

Solution of the Problem

We define Laplace transform and Fourier transform respectively as:

$$\bar{f}(s,x_1,x_3) = \int_0^\infty f(t,x_1,x_3)e^{-st}dt\ , \qquad (25)$$

$$\hat{f}(x_3,\xi,s) = \int_{-\infty}^{\infty}\bar{f}(s,x_1,x_3)e^{\iota\xi x_1}dx_1\ , \qquad (26)$$

Applying Laplace transform defined by (25) on (20)-(24) and then applying Fourier transforms defined by (26) on the resulting quantities and eliminating$\hat{C}\&\hat{T}$, $\hat{\phi}\&\hat{T}$,$\hat{\phi}\&\hat{C}$ and $\hat{\phi}_2$ respectively from the resulting equations, we obtain:

$$[D^6 + AD^4 + BD^2 + C]\hat{\phi} = f_1 e^{-\gamma^* x_3} \qquad (27)$$

$$[D^6 + AD^4 + BD^2 + C]\hat{T} = f_2 e^{-\gamma^* x_3} \qquad (28)$$

$$[D^6 + AD^4 + BD^2 + C]\hat{C} = f_3 e^{-\gamma^* x_3} \qquad (29)$$

$$[D^4 + ED^2 + F]\hat{\psi} = 0\ , \qquad (30)$$

The solutions of the equations (27)-(30) satisfying the radiation conditions that $\left(\hat{\phi},\widehat{\phi^*},\hat{T},\widehat{\phi_2},\hat{\psi}\right) \to 0$ as $x_3 \to \infty$ are given by:

$$\hat{\phi} = B_1 e^{-m_1x_3} + B_2 e^{-m_2x_3} + B_3 e^{-m_3x_3} + L_1 e^{-\gamma^* x_3} \quad (31)$$

$$\hat{T} = d_1B_1 e^{-m_1x_3} + d_2B_2 e^{-m_2x_3} + d_3B_3 e^{-m_3x_3} +_2 e^{-\gamma^* x_3} \qquad (32)$$

$$\hat{C} = e_1B_1 e^{-m_1x_3} + e_2B_2 e^{-m_2x_3} + e_3B_3 e^{-m_3x_3} + L_3 e^{-\gamma^* x_3} \qquad (33)$$

$$\hat{\psi} = B_4 e^{-m_4x_3} + B_5 e^{-m_5x_3} \qquad (34)$$

$$\widehat{\phi_2} = h_4B_4 e^{-m_4x_3} + h_5B_5 e^{-m_5x_3} \qquad (35)$$

where $d_i = \frac{a_{39}m_i^4 - a_{40}m_i^2 + a_{41}}{a_{37}m_i^2 + a_{38}}, e_i = \frac{a_{42}m_i^4 + a_{48}m_i^2 + a_{44}}{a_{37}m_i^2 + a_{38}}\ ,\ i = 1,2,3\& h_i = \frac{a_2(m_i^2 - \xi_1)}{a_3}\ , i = 5,6$

$L_i = \frac{f_i}{[m_i^6 + Am_i^4 + Bm_i^2 + C]}, i = 1,2,3$

and $m_i^2 (i = 1,2,3)$ are the roots of the characteristic equation of equation (27) and $m_i^2 (i = 4,5)$ are the roots of the characteristic equation of equation (30).

Boundary Condition

We consider concentrated normal force and concentrated thermal source at the boundary surface$x_3 = 0$, mathematically, these can be written as:

$$t_{33} = -F_1\psi_1(x_1)\delta(t), t_{31} = 0, m_{32} = 0,$$
$$T = F_2\psi_1(x_1)\delta(t)\ ,\ C = F_3\psi_1(x_1)\delta(t) \qquad (36)$$

Where F_1 is the magnitude of the applied force and F_2 is the constant temperature applied on the boundary.
Also

$$t_{33} = \lambda e + (2\mu + K)u_{3,3} - \beta_1\left(1 + \tau_1\frac{\partial}{\partial t}\right)T - \beta_2\left(1 + \tau^1\frac{\partial}{\partial t}\right)C$$

$t_{31} = (2\mu + K)u_{3,1} - K\phi_2 m_{32} = \beta\phi_{2,3}$ (37)

Substituting the values of $\hat{\phi}, \widehat{\phi^*}, \hat{T}, \hat{\psi}, \widehat{\phi_2}$ from the equations (31)-(35) in the boundary condition (36) and using (5)-(11), (12)-(13), (25)-(26) and solving the resulting equations, we obtain:

$\widehat{t_{33}} = \sum_{i=1}^{5} G_{1i} e^{-m_i x_3} - M_1 e^{-\gamma^* x_3}, i = 1,2, \dots, 5$ (38)

$\widehat{t_{31}} = \sum_{i=1}^{5} G_{2i} e^{-m_i x_3} - M_2 e^{-\gamma^* x_3}, i = 1,2, \dots, 5$ (39)

$\widehat{m_{32}} = \sum_{i=1}^{5} G_{3i} e^{-m_i x_3} - M_3 e^{-\gamma^* x_3}, i = 1,2, \dots, 5$ (40)

$\hat{T} = \sum_{i=1}^{5} G_{4i} e^{-m_i x_3} - M_4 e^{-\gamma^* x_3}, i = 1,2, \dots, 5$ (41)

$\sum_{i=1}^{5} G_{5i} e^{-m_i x_3} - M_5 e^{-\gamma^* x_3}, i = 1,2, \dots, 5$ (42)

Case 1: for the thermal source: $F_1 = 0$

Case 2: for the normal source: $F_2 = 0$

Applications

Uniformly distributed source: The solution due to uniformly distributed force applied on the half-space is obtained by setting

$$\psi_1(x_1) = \begin{cases} 1, & |x_1| \le d \\ 0, & |x_1| > d \end{cases} \quad (43)$$

Applying Laplace and Fourier transforms on (4.7), gives $\widehat{\psi_1}(\xi) = \frac{2\sin(\xi d)}{\xi}, \xi \neq 0$ (44)

Linearly distributed source

The solution due to linearly distributed force over a strip of non-dimensional width 2d, applied on the half-space is obtained by setting

$$\psi_1(x_1) = \begin{cases} 1 - \frac{|x_1|}{d}, & |x_1| \le d \\ 0, & |x_1| > d \end{cases} \quad (45)$$

Applying Laplace and Fourier transforms on (4.7), gives

$$\widehat{\psi_1}(\xi) = \frac{2(1-\cos(\xi d))}{\xi^2 d}, \xi \neq 0 \quad (46)$$

Particular cases

(i) If we take $\tau_1 = \tau^1 = 0,\ \varepsilon = 1$, in Eqs. (38)- (42), we obtain the corresponding expressions of stresses, displacements and temperature distribution for microstretch thermoelastic half space with one relaxation time.

(ii) If we take $\varepsilon = 0$ in Eqs. (38)- (42), the corresponding expressions of stresses, displacements and temperature distribution are obtained for microstretch thermoelastic half space with two relaxation times.

(iii) Taking $\tau^0 = \tau^1 = \tau_0 = \tau_1 = 0$ in Eqs. (38)- (42), yield the corresponding expressions of stresses, displacements and temperature distribution for microstretch coupled thermoelastic half space.

Special case

Micropolar Thermoelastic Solid

In absence of mass diffusion effect in Equations (38) - (42), we obtain the corresponding expressions of stresses, displacements and temperature for micropolar generalized thermoelastic half space.

Inversion of the transforms

The transformed displacements, stresses and temperature changes are functions of the parameters of Laplace and Fourier transforms s and ξ respectively and hence these are of the form $f(s, \xi, z)$. To obtain the solution of the problem in the physical domain, we must invert the Laplace and Fourier transform by using the method applied by Kumar [34].

Numerical Results and Discussion

The analysis is conducted for a magnesium crystal-like material. For numerical computations, following Eringen [3],the values of physical constants are:

$\lambda = 9.4 \times 10^{10} \mathrm{Nm}^{-2}, \mu = 4.0 \times 10^{10} \mathrm{Nm}^{-2}, K = 1.0 \times 10^{16} \mathrm{Nm}^{-2}, \rho = 1.74 \times 10^{3} \mathrm{Kgm}^{-3}, j = 0.2 \times 10^{-19} \mathrm{m}^2, \quad \gamma = 0.779 \times 10^{-9} \mathrm{N}$

Comparison of the dimensionless form of the field variables for the cases of micropolar mass diffusion thermoelastic medium (MPMD) and micropoar thermoelastic medium (MP) for two different values of time t (*t=.01* and *t=.02*), subjected to linearly distributed source is shown in Figures 1-5. The values of all physical quantities for all cases are shown in the range $0 \le x_1 \le 2$.Solid lines, dash lines corresponds to micropolar thermoelastic mass diffusion medium (MPMDT1) for t *=0.01*andmicropolar thermoelastic mass diffusion medium (MPMDT2)for*t=0.02* respectively. Solid lines with central symbol & dash line with central symbol corresponds to micropolar thermoelastic (MPT1 and MPT2) for *t=.01*and *t=.02*respectively.

Linearly distributed normal force:

Fig. 1 shows the variation of normal stress t_{33} with the distance x_1. It is noticed that for MPMDT1 and MPMDT2, t_{33} show similar behavior. The value of normal stress monotonically increases as x_1 and then oscillates. The value of t_{33} increases near the application of the normal force due to the mass diffusion effect and then remain oscillating for all values of x_1.

Fig. 2 displays the variation of tangential stress t_{31} with the distancex_1. It is noticed that initially the behavior of t_{31} for MPMDT1 and MPT1 show variable trend but for MPMDT1, MPMDT2 and MPT1, MPT2 exhibits similar behavior. t_{31} Initially decrease monotonically for all the cases. The variation in tangential stress in micropolar thermoelastic is more than that of micropolar thermoelastic with mass diffusion.

Fig. 3 shows the variation of couple tangential stress m_{32} with distancex_1. The behavior and variation of $m_{32}$for MPMDT1, MPMDT2, MPT1 and MPT2 remain similar to each other for all values of x_1. The magnitude of couple tangential stress in micropolar thermoelastic with mass diffusion is more than that of micropolar thermoelastic

Fig. 4 depicts the variation of temperature$T$ with distance x_1. The trend and variation of T is similar in case of MPMDT1, MPMDT2 and MPT1 initially. For these curves the initial behavior is monotonically decreasing and oscillator away from the point of application of normal force. MPT2 show opposite trend initially. Fig. 5 display the variation of mass concentration$C$ with distancex_1. For MPMDT1 and MPMDT2 the graphs are similar. Initially the trend is decreasing. After some oscillatory behavior mass concentration approaches to the boundary surface away from the application of force.

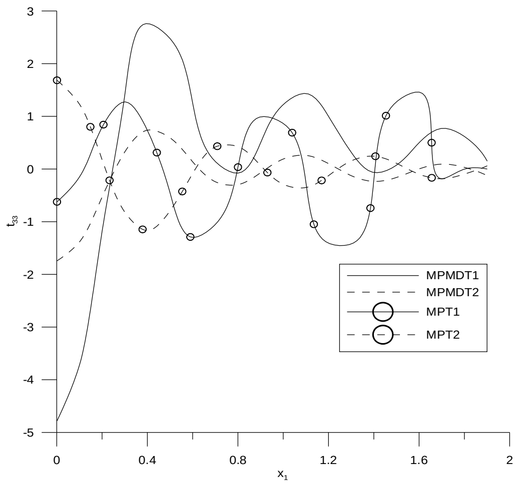


Fig. 1. Variation of normal stress withx_1,

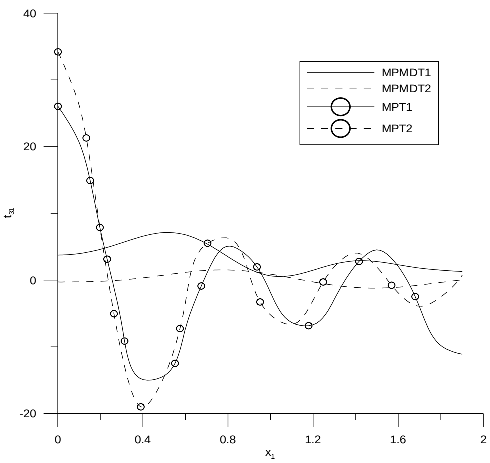
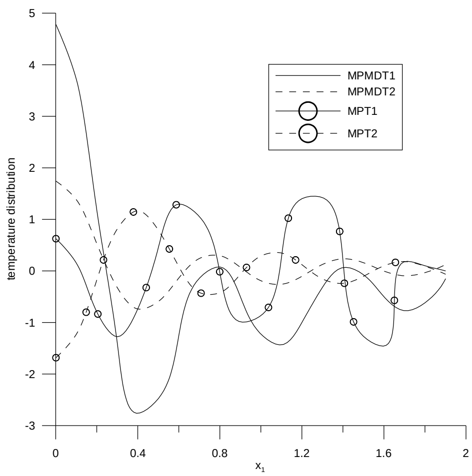


Fig. 2. Variation of tangential stress withx_1



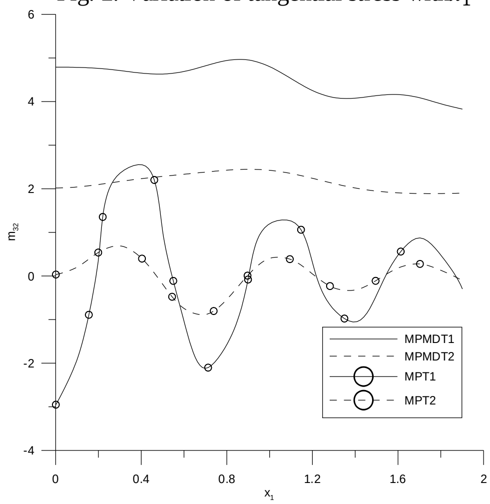


Fig. 3. Variation of couple tangential stress with x_1 Fig. 4. Variation of temperature distribution with x_1

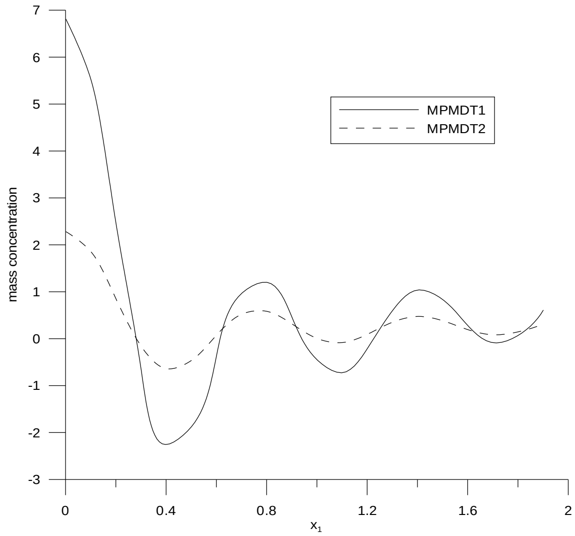


Fig. 5. Variation of mass concentration with x_1

Conclusions

The problem consists of investigating displacement components, scalar mass concentration, temperature distribution and stress components in a homogeneous isotropic micropolar mass diffusion thermoelastic half space due to various sources subjected to laser pulse. Integral transform technique is employed to express the results mathematically. Theoretically obtained field variables are also exemplified through a specific model to present the results in the transformed domain. The analysis of results permits some concluding remarks:

(1) It is clear from the figures that all the field variables have nonzero values only in the bounded region of space indicating that all the results are in agreement with the generalized theory of thermoelasticity.

(2) The effect of the mass concentration is much pronounced in all the resulting quantities.

(3) It is noticed that the figures that the time t plays a significant role in all the field quantities. Changes inthe value of time t cause significant changes in all the simulated resulting quantities.

(4) It can be easily concluded from the figures that the curves for various stresses in case of micropolar mass diffusion solid show similar trends.

(5) The variation of mass concentration differs significantly due to the presence of normal force and due to the presence of thermal source.

(6) Tangential stress, couple stress and temperature change are also affected due to diffusion effect as well as load/source applied.

The new model is employed in a micropolar mass diffusion thermoelastic medium as a new improvement in the field of thermoelasticity. The subject becomes more interesting due to irradiation of a laser pulse with an extensive short duration or a very high heat flux has found numerous applications. The method used in this article is applicable to a wide range of problems in thermodynamics. By the obtained results, it is expected that the present model of equations will serve as more realistic and will provide motivation to investigate microstretch generalized thermoelasticity problems regarding laser pulse heat with high heat flux and/or short time duration.

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